First-principles Investigation on Structural and Electronic Properties of Cubic Na₂O

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Abstract

The structural and electronic properties of cubic Na₂O were investigated by first-principles calculations. The calculated structural parameters $(a, B_0, \text{ and } B'_0)$ show a good agreement with the available values. Furthermore, the electronic band structure and density of states were obtained. We found that the cubic Na₂O is a direct band gap material. In addition, the electron localized function (ELF) was analyzed to give more explanation of the bonding nature of cubic Na₂O. The ELF analysis indicates that the Na-O band is a typical ionic bond. The analysis details were presented and discussed.

Keywords: Na₂O, electronic structure, first-principles

1. Introduction

It is well known that Na₂O is the alkali metal oxide and crystallizes in the cubic structure (antifluorite type, space group $Fm\overline{3}m$). Na₂O has been reported to be a promising candidate for various technological applications such as solid-state batteries, fuel cells, and solid-state gas detectors [1-3]. It is difficult to prepare and preserve Na₂O at ambient condition. Thus, a little experimental study has been reported for Na₂O. Zintl et al. [4] determined the room temperature lattice parameter of Na₂O by using a powder diffraction experiment. Recently, Wu *et al.* [5] studied the high-pressure behavior of Na_2O up to 30 GPa by synchrotron angledispersive powder X-ray diffraction in a diamond anvil cell at room temperature. On the theoretical side, the lattice constants, and elastic properties of Li₂O, Na₂O, and K₂O were computed via the ab initio Hartree-Fock LCAO method by Dovesi et al. [6]. This method also was used to study the stability of the alkali metal oxides under pressure by Cančarević et al. [7]. Moakafi et al. [8] reported results of first-principles calculations for the electronic and optical properties under pressure effect of Li₂O, Na₂O, Ki₂O, and Rb₂O compounds in the cubic antifluorite structure. Thompson et al. [9] presented a theoretical study of electronic and phonon spectra of Na₂O using density functional theory with a plane-wave in both localdensity approximation and generalized gradient approximation. In this work, the firstprinciples calculations with the generalized gradient approximation (GGA) of Perdew-Burke-Ernzerhof (PBE) were carried out to investigate the structural and electronic properties of cubic Na₂O. The structural parameters, electronic band structure, density of states (DOS), and electron localized function (ELF) of cubic Na₂O were presented and discussed.

2. Computational details

First-principles calculations were performed using Vienna Ab-inito Simulation Package (VASP) [10]. The generalized gradient approximation (GGA) proposed by Perdew-Burke-Ernzerhof (PBE) was used to describe an exchange-correlation functional [11-12].



Fig. 1. Crystal structure of cubic Na₂O. The Na and O atoms are represented with gray and red balls, respectively.

The projector-augmented-wave (PAW) method was treated to perform the electron wave functions [13]. For the plane-wave expansion of the wave functions, a cut off energy of 600 eV was used. The Brillouin zone integrations were performed using Monkhorst-pack grids with $8 \times 8 \times 8$ *k*-points mesh to obtain well convergence of total energy [14]. The total change in energy less than 10^{-4} eV and the force less than 10^{-3} eV/Å were used as convergence criteria for structural relaxation. The unit cell of cubic Na₂O used in this calculation consists of 12 atoms per unit cell (Na 8 atoms and O 4 atoms). The cubic structure of Na₂O is shown in Fig.1.

3. Results and discussion

3.1 Structural properties

The total energies as a function of volume, E (V), of a unit cell of cubic Na₂O were calculated and fitted into the Birch-Murnaghan's equation of state [15-17] to obtain the equilibrium volume (V_0), the bulk modulus (B_0), and its pressure derivatives (B'_0). The Birch-Murnaghan's equation of state is written as,

$$E(V) = E_0 + \frac{9}{16}V_0B_0 \left\{ \left[\left(\frac{V_0}{V}\right)^2 - 1 \right]^3 B'_0 + \left[\left(\frac{V_0}{V}\right)^2 - 1 \right]^2 \left[6 - 4\left(\frac{V_0}{V}\right)^2 \right] \right\},\tag{1}$$

where E_0 is the equilibrium energy, V_0 is the equilibrium volume, B_0 is the bulk modulus, and B'_0 is its pressure derivative. The relationship between total energy and volume (*E-V*) of cubic Na₂O is shown in Fig. 2. The *E-V* data was fitted to the Birch-Murnaghan's equation of state and the fitting parameters ($a = \sqrt[3]{V_0}$, B_0 , and B'_0) are listed in Table 1. Our calculated results agree with previous calculations. Note that our calculated lattice constant is larger than the experimental value, reflecting in a smaller in the bulk modulus. However, the disagreement between our calculated lattice constant and the experimental value is only 1.95%, indicating that the calculations are reliable.



Fig. 2. The relationship between total energy and volume of cubic Na₂O. The solid line represents Birch-Murnaghan equation of state fits to the calculated data (solid points).

Table 1. Calculated lattice constant (*a*), bulk modulus (*B*), pressure derivative of bulk modulus (B'_0) along with the previous calculations and available experimental data of cubic Na₂O.

| | a (Å) | B (GPa) | B'_0 |
|--------------|-----------------|-----------------|----------|
| This work | 5.597 | 45.40 | 4.35 |
| Calculations | 5.398(LDA) [9] | 56.00 (LDA) [9] | |
| | 5.408 (LDA) [8] | 57.79 (LDA) [8] | 4.20 [8] |
| | 5.583 (GGA) [9] | 54.00 (GGA) [9] | 4.71 [8] |
| | 5.592 (GGA) [8] | 47.11 (GGA) [8] | |
| Experiment | 5.490 [4] | 55.60 [5] | - |

3.2 Electronic properties

The calculated energy band structure of cubic Na₂O along with the high symmetry point of the Brillouin zone is shown in Fig. 3. The energy band gap of cubic Na₂O was calculated to be 1.88 eV with a direct gap, which is in good agreement with the previous GGA calculation value of 1.83 eV from Moakafi *et al.* [8]. To further explain the electronic band structure, the total and partial density of states (DOS) also were calculated as shown in Fig. 4. We found that the main bonding peaks located close to the Fermi level (or the upper valence band (VB)) at the energy range between -3 and 0 eV. The contribution of O-p, Na-s, and Na-p states are dominant. However, the O-p state is mainly dominant at the Fermi level. For investigating the bonding nature of materials, the electron localization function (ELF) is used to analyzes the characteristics and types of chemical bonds [18]. The ELF = 1 shows the strongest covalent bond, ELF = 0.5 shows the metallic bond, and $0 \le \text{ELF} < 0.5$ shows the ionic bond. The electron localization function (ELF) of cubic Na₂O was calculated and shown in Fig. 5. It is clearly seen that the Na-O bonds are in the blue area which corresponds to the ELF < 0.5. Thus, the bond formed between Na and O is an ionic bond.



Fig. 3. Calculated band structure of cubic Na₂O.



Fig.4. Calculated total and partial density of states (DOS) of cubic Na₂O.



Fig. 5. Calculated electronic localization functions (ELF) of cubic Na_2O . Electron localization function isosurface maps for (a) Miller indices: (0 0 1) and (b) Miller indices: (1 1 0).

4. Conclusion

The first-principles calculations based on density functional theory within PBE-GGA were used to investigate the structural and electronic properties of cubic Na₂O. The calculated structural parameters such as lattice constant, bulk modulus, and its pressure derivative are in good agreement with the available data. The electronic band structure and the density of states were determined. The cubic Na₂O was found to be a direct band gap material with the band gap value of 1.88 eV, which is consistent with the previous study. Moreover, the electron localized function (ELF) was also analyzed and the result indicates that the Na-O is ionic bonding. We hope that this work will be used as a reference data for both theoretical and experimental studies based on Na₂O compound.

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References

- Jamal M, Venugopal G, Shareefuddin M, Chary MN. Sodium ion conducting glasses with mixed glass formers NaI–Na₂O–V₂O₅–B₂O₃: application to solid state battery. Materials Letters. 1999;39(1):28-32.
- [2] Cho KI, Lee SH, Cho KH, Shin DW, Sun YK. Li₂O–B₂O₃–P₂O₅ solid electrolyte for thin film batteries. J. Power Sources. 2006;163(1):223-8.
- [3] Lee KB, Beaver MG, Caram HS, Sircar S. Performance of Na₂O promoted alumina as CO₂ chemisorbent in sorption-enhanced reaction process for simultaneous production of fuel-cell grade H₂ and compressed CO₂ from synthesis gas. J. Power Sources. 2008; 176(1):312-9.
- [4] Zintl E, Harder A, Dauth B. Lattice Structure of the oxides, sulfides, selenides and tellurides of lithium, sodium and potassium. Z Elektrochem. 1934;40:588.
- [5] Wu X, Zhang Y, Zhang J, Liu R, Yang J, Yang B, Xu H, Y Ma. High pressure X-ray diffraction study of sodium oxide (Na₂O): Observations of amorphization and equation of state measurements to 15.9 GPa. J. Alloys Compd. 2020;823:153793.
- [6] Dovesi R, Roetti C, Freyria-Fara C, Prencipe M, Saunders VR. On the elastic properties of lithium, sodium and potassium oxide. An ab initio study. Chem. Phys. 1991;156:11-9.
- [7] Čančarević Z, Schön JC, Jansen M. Stability of alkali-metal oxides as a function of pressure: Theoretical calculations. Phys. Rev. B: Condens. Matter. 2006;73:224114.
- [8] Moakafi M, Khenata R, Bouhemadou A, Khachai H, Amrani B, Rached D, Rérat M. Electronic and optical properties under pressure effect of alkali metal oxides. Eur. Phys. J. B. 2008;64(1):35-42.
- [9] Thompson M, Shen X, Allen PB. Density functional calculation of electronic structure and phonon spectra of Na₂O. Phys. Rev. B. 2009;79(11):113108.
- [10] Kresse G, Furthmuller J. Efficient iterative schemes for ab initio total-energy calculations using a plane-wave basis set. Phys. Rev. B. 1996;54:11169-86.
- [11] Perdew JP, Chevary JA, Vosko SH, Jackson KA, Pederson MR, Singh DJ, Fiolhais C. Atoms, molecules, solids, and surfaces: Applications of the generalized gradient approximation for exchange and correlation. Phys. Rev. B: Condens. Matter. 1992; 46(11):6671-87.
- [12] Perdew JP, Burke K, Ernzerhof M. Generalized Gradient Approximation Made Simple. Phys Rev Lett. 1996;77(18):3865-8.

- [13] Blöchl P E. Projector augmented-wave method, Phys. Rev. B: Condens. Matter. 1994;50: 17953–79.
- [14] Monkhorst J, Pack JD. Special points for Brillouin-zone integrations, Phys. Rev. B: Condens. Matter. 1976; 13:5188 –92.
- [15] Birch F. Finite Elastic Strain of Cubic Crystals. Phys. Rev. B: Condens. Matter. 1947; 71(11):809-24.
- [16] Murnaghan FD. Finite Deformations of an Elastic Solid. Am. J. Math. 1937;49:235-60.
- [17] Murnaghan FD. The Compressibility of Media under Extreme Pressures. Proc Natl Acad Sci U S A. 1944;30:244-7.
- [18] Burdett JK, Mc Cormick TA. Electron localization in molecules and solids: The meaning of ELF. J. Phys. Chem. A 1998;102:6366-72.